

1 Detrended Fluctuation Analysis of Spatially Extended Digital Surfaces:
2
3 the Classification Process of $1/f$ Noise and Computacional Performance

4 Vanessa Cristina Oliveira de Souza^{a, b1} and Arcilan Trevenzolli Assireu^b
5 ^aNational Institute for Space Research, São José dos Campos, SP, Brazil
6 ^bFederal University of Itajubá, Itajubá, MG, Brazil

7 *Received on Nov 02, 2016 / accepted on Dec 29, 2016*

8

9

Abstract

10 A two-dimensional Detrended Fluctuation Analysis (DFA-2D) was
11 employed for the characterization of $\frac{1}{f}$ classic white, pink and brownian
12 noises. The values were compared with the fBm and fGn fractal signals.
13 The exponent α was evaluated and compared with exponent β of the
14 power spectrum. It was determined that the DFA-2D was capable
15 of detecting and characterizing the different signals. In addition, the
16 relation between α and β was defined as $\beta \equiv 2\alpha - 2$ for two-dimensional
17 signals, with the exception of the fBm signals, where $\beta \equiv 2\alpha$. A
18 computational analysis of the time and accuracy of the DFA-2D was
19 also carried out. The complexity of the method grows exponentially
20 with the enlargement of the image size and the process becomes very
21 slow for images greater than 1000x1000 pixels.

22

23 **Keywords:** long-range dependence, DFA-2D, $\frac{1}{f}$ noise, computational
24 analysis.

25

26 1 Introduction

27 The method employed in the Detrended Fluctuation Analysis (DFA) was
28 formally defined by Peng et al [1] for an analysis of non-stationary signals.
29 In 2002, the method was generalized to characterize multifractal signals[2].
30 More recently, the DFA and multifractal DFA (MF-DFA) have been ex-
31 panded to operate not only for a one-dimensional time series but also for
32 any other kind of dimensions, in particular two-dimensional surfaces (DFA-
33 2D and MF-DFA2D)[3]. As well as this, there have been several other
34 suggestions for generalizing the method in recent years [4].

¹E-mail Corresponding Author: vanessasouza@unifei.edu.br

35 DFA is regarded as a method that is able to capture the behavior of
36 scales (power law) in long-range correlations in time series, which, in the
37 analysis of signals, is a concept also known as memory or persistence. The
38 DFA' differential is the fact that the method operates in a non-stationary
39 series. If the signal is detrended, it is able to calculate the correlations
40 without being affected by the possible presence of trends. Moreover, as the
41 DFA measures long-range correlations, it is also suitable for calculating what
42 is called the Hurst Exponent [5]. In this way, the fractal theory forms the
43 underlying framework of the DFA.

44 Although some authors have criticized the method [6, 7, 8, 9], the one-
45 dimensional DFA has been widely used, for example in biomedical [10, 11, 12]
46 and environmental areas [13, 14]. Its multifractal version is also popular,
47 because many of the dynamic processes that occur in nature, cannot be
48 characterized by a single scaling exponent [15, 16, 17].

49 When applied to two-dimensional surfaces, the notion of persistence can
50 reveal features related to surface roughness. For this reason, DFA-2D and
51 MF-DFA2D have been particularly applied in these conditions. However,
52 although it is an established method for its one-dimensional version, the
53 DFA-2D has not been explored to any extent. Some of its applications can
54 be seen in [18, 19, 20, 21, 22, 23, 24].

55 One important factor that should be mentioned is that some authors use
56 a one-dimensional DFA to operate within a two-dimensional matrix [25, 26,
57 27]. The main reasons given by these authors is the high computational cost
58 of the DFA-2D and the difficulty of interpreting the results.

59 Most of the works in the literature only use the DFA-2D as a computer
60 operator, without being particularly concerned about the physical interpre-
61 tation of the results. In addition, theoretical studies in general are carried
62 out with knowingly defined fractal data such as the Fractional Brownian
63 Motion (fBm) and the Fractional Gaussian Noise (fGn).

64 In light of this, this study seeks to assist in the understanding and inter-
65 pretability of the DFA-2D, by examining the two-dimensional DFA through
66 a canonical data model that includes both fractal and $\frac{1}{f}$ noises and the re-
67 lation between the exponent of the DFA-2D (α), with that of the 2D Power
68 Spectrum (β). Another aim of this work was to assess the computacional
69 performance of DFA-2D with different image sizes and different Hurst ex-
70 ponents.

71 This study is structured in the following way: Section 2 provides a short
72 review of the data used in the research. Section 3 describes the DFA-2D
73 algorithm and the data classification process, as well as conducting a com-
74 putational analysis. Section 4 concludes the study and adds some final

75 considerations.

76

77 **2 $\frac{1}{f}$ Noise Data**

78 [28] states that several physical systems have some type of power-law cor-
79 relations in space (fractal) or time ($\frac{1}{f^\gamma}$). [29] define 1/f noise as a kind of
80 temporal fluctuation which has an intensity and power that are inversely
81 proportional to frequency. The noises of the 1/f type vary their correlation
82 power between the white ($\frac{1}{f^0}$) and the brownian noise($\frac{1}{f^2}$). There are several
83 other noises between these two extremes, including the intermediary pink
84 noise ($\frac{1}{f}$), which contains influences that are equal to all the intervals of the
85 time scale in the whole spectrum [30].

86 In contrast with the white noise that does not have memory, the brown-
87 ian processes are dominated by their recent history. Both are non-stationary.
88 The pink noise shares properties with the white and brownian noises and
89 is considered to be non-stationary and as a result of these intermediary
90 features is the most widely used model to design real events.

91 When in two-dimensional space, these concepts can be extended to spa-
92 tial autocorrelation, which is defined as the correlation between pairs of
93 separate points by use of distance. [31] explains that the fractals have a
94 form that is simple but important for the spatial structure, since their pat-
95 terns are statistically equivalent in all the scales. Thus, the $\frac{1}{f^\gamma}$ family of
96 noise models can be regarded as fractional brownian motion, with the ex-
97 ception of $\gamma \equiv 1$ [30]. In this case the range of possible values for γ is altered.
98 Brownian fractional motion are defined by low frequencies, with $1 < \gamma < 3$,
99 including the ordinary brownian motion ($\gamma = 2$). However, the processes
100 derived from fractional gaussian noises are defined by high frequencies, with
101 $-1 < \gamma < 1$, including the white gaussian noise ($\gamma = 0$) [32].

102 The data regarding white gaussian, pink and brownian noises used in
103 this study, were generated through the Matlab code by [33], based on the
104 methodology described in [31], and using the Inverse Fourier Transform.
105 The size of the matrices was 500x500. Five hundred samples of each model
106 were produced.

107 Figure 1 shows a sample of each model and its respective power spectrum.
108 Since it is two-dimensional, the exponent β is calculated by means of a radial
109 average and in this case, it expresses the average of the 2D Fourier power
110 spectrum in all possible directions [34]. The Matlab code distributed by [35]
111 was used for power spectrum computation.

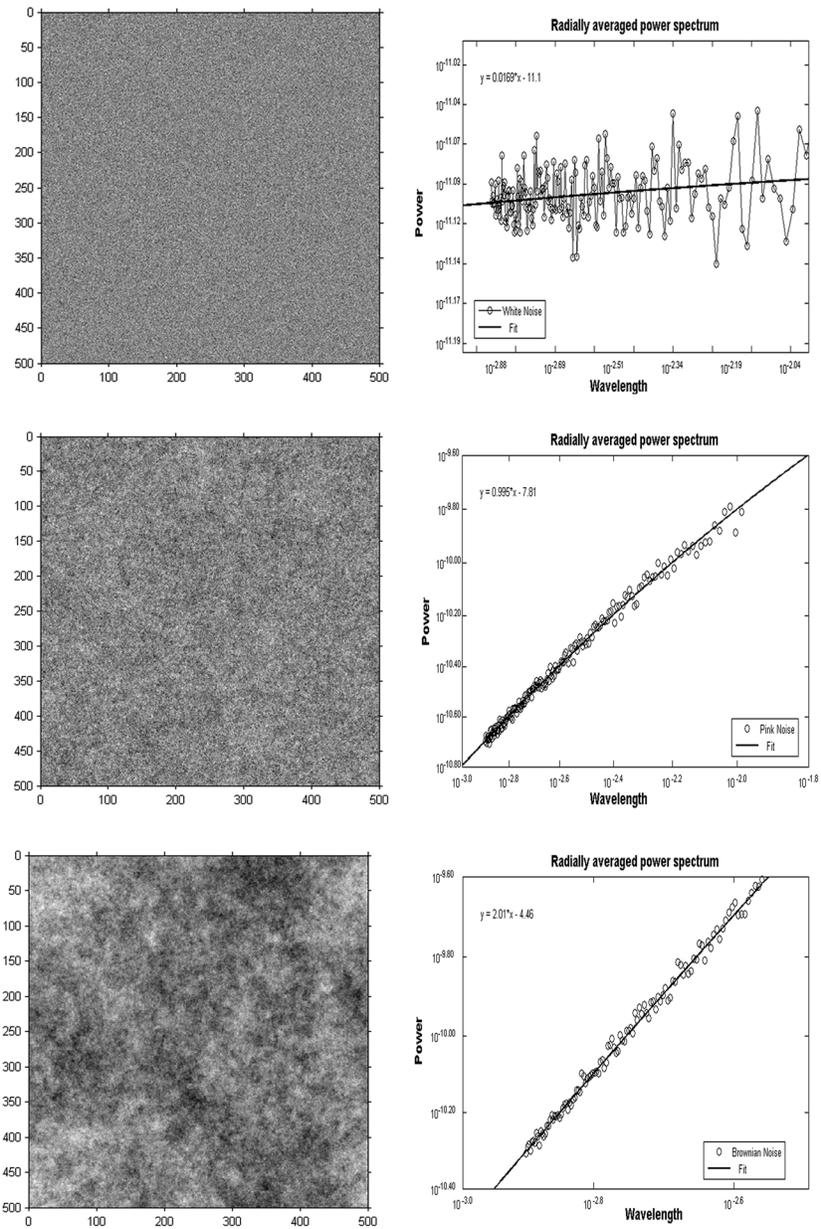


Figure 1: Examples of noisy surfaces and their respective radial power spectrum. At the top, white noise. In the middle, pink noise and down, Brownian noise.

Signal	β_{mean}
White Noise	-0.0009 ± 0.0083
Pink Noise	1.0022 ± 0.0080
Brownian Noise	2.0051 ± 0.0074

Table 1: Power spectrum values for the radially averaged power spectral density for some typical signals. The mean and standard deviation are relative to five hundred samples of each model.

112 Table 1 shows the average values and respective deviations of the power
 113 spectrum. It was confirmed that the values obtained are identical to those
 114 usually described in the literature for the one-dimensional time series. In
 115 the case of simulations of fBm data, when the Hurst exponent = 0.5, $\beta =$
 116 2.7597 ± 0.094 was obtained. As for fGn signals also with Hurst exponent
 117 = 0.5, $\beta = -0.9839 \pm 0.0134$ was obtained.

118

119 3 Detrended Fluctuation Analysis 2D (DFA-2D)

120 3.1 The DFA-2D Algorithm

121 The implementation of the DFA-2D method is an adaptation of the DFA
 122 method for a one-dimensional temporal series and will be described below,
 123 as defined in [3].

124 In the 2D adaptation, the matrix is subdivided into submatrices that
 125 range from 6 to $(n/4)$, n being the lowest value between $(M/4)$ e $(N/4)$,
 126 where M and N correspond to the number of lines and columns of the
 127 matrix respectively. After the subdivision of the matrix, it is integrated in
 128 accordance with Equation 1.

$$y_{v,w}(i, j) = \sum_{k_1=1}^i \sum_{k_2=1}^j X_{v,w}(k_1, k_2) \quad (1)$$

129 The local trend is removed for each sub-matrix. The least-squares method
 130 is employed by adjusting the surface by one of the surface equations defined
 131 in [3]. Following this, the residual matrix between the sub-matrix (y) and
 132 the adjusted surface (\tilde{y}) is calculated (Equation 2).

$$z_{v,w}(i, j) = y_{v,w}(i, j) - \tilde{y}_{v,w}(i, j) \quad (2)$$

133 The local fluctuation function for each scale s ($6 \leq s \leq \min((M, N)/4)$)
 134 is calculated by means of Equation 3.

$$F^2(u, v, s) = \frac{1}{s^2} \sum_{i=1}^s \sum_{j=1}^s (z_{v,w}(i, j))^2 \quad (3)$$

135 The global fluctuation function is calculated as an average of the local
 136 fluctuation functions, by means of Equation 4.

$$F_2(s) = \left(\frac{1}{p * q} \sum_{i=1}^p \sum_{j=1}^q F^2(u, v, s) \right)^{\frac{1}{2}} \quad (4)$$

137 where: $p = M/s$ and $q = N/s$.

138 Thus by varying the value of s , it is possible to determine the relation of
 139 the scale between the fluctuation of the $F_2(s)$ function and the size of Scale
 140 s , by means of Equation 5:

$$F(s) \sim s^\alpha \quad (5)$$

141 The exponent for fluctuation scaling (α) is thus calculated as the slope
 142 of the straight line of the log-log graph between $F_2(s)$ and s .

143 3.2 The exponent α

144 According to some authors, one of the drawbacks of the DFA is the inter-
 145 pretability of exponent α . Since the DFA is used as an approximation for
 146 the Hurst exponent, the interpretation in the literature generally follows
 147 the pattern of Table 2, which, in reality, describes the possible values and
 148 interpretations of the Hurst exponent. The exponent α does not only range
 149 between 0 and 1. In general terms [36, 37]:

- 150 • If the signal is a fractional gaussian noise (fGn), i.e. is stationary, then
 151 $H \equiv \alpha$.
- 152 • If the signal is a fractional brownian motion (fBm), i.e. is not station-
 153 ary, then $H \equiv \alpha - D$, where D is the fractal dimension of the analyzed
 154 data.

155 This relationship arises from the first stage of the DFA algorithm, relative
 156 to the integration of the signal. The DFA acts on the fBm signals and
 157 the integration turns a possible fGn signal into a fBm and a fBm into an
 158 accumulated fBm [37, 38].

α	Interpretation
$0 < \alpha < 0.5$	Signal with long-range anti-persistence correlation. The closer to zero the higher the anti-persistence.
$\alpha = 0.5$	Uncorrected signal (white noise)
$0.5 < \alpha < 1$	Signal with persistent long-range correlation. The closer to 1, the greater the persistence.
$\alpha = 1$	Pink noise
$\alpha = 1.5$	Brownian noise

Table 2: Values and interpretation of the exponent α

159 Hence, the stationary signals return values of α between 0 and 1 for the
 160 DFA-2D, whereas the non-stationary signals return values between 2 and
 161 3. Values of α between 1 and 2 indicate that the signal does not possess
 162 memory. An understanding of the significance of the value of α helps to
 163 understand not only the correlation of the signal power analyzed, but also
 164 its origins.

165 In the case of a one-dimensional series, [39] defined the power spectrum
 166 β [40] as the equivalent of $\beta = 2H + 1$, for fBm signals and $\beta = 2H - 1$
 167 for fGn signals. [41, 39] provide an experimental validation of the relation
 168 between β and α , both for fBm, and for fGn, as:

$$\beta \equiv 2\alpha - 1 \quad (6)$$

169 3.3 Implementation of the DFA-2D

170 The code was designed with the aid of Python language and in particular the
 171 NumPy and SciPy [42]. Python version 3.4, was used together with numPy
 172 1.11 and sciPy 0.17. NumPy does not have direct functions for calculating
 173 the two-dimensional least squares in the matrix. In view of this, a LU
 174 factorization was employed to carry out the detrending of each sub-matrix
 175 on the basis of the plane model $\tilde{y}_{v,w}(i, j) = ai + bj + c$.

176 The DFA-2D was evaluated in scales ranging from 6 to 128. The empiri-
 177 cal tests were carried out on a desktop, with a Core i7 processor 950, 3.07
 178 GHZ, 8GB of RAM memory, Linux Centos logrotate 7.2.15.11, with totally
 179 dedicated servers.

180 The experiment conducted by [3] was reproduced as it is, to validate
 181 if the method had been implemented correctly i.e. by varying the Hurst
 182 exponent between 0.05 0.95, with increments of 0.05. 500 simulations of fBm
 183 fractals were produced for each value of the Hurst exponent with dimensions
 184 of 500x500, and using Fraclab [43]. The graph in Figure 2 shows the results.

185 The error bar shows the standard deviation of the 500 samples of each value
 186 of H . Thus as stated by [3], the standard deviation increases with the increase
 187 of H .

188 The Python code is available at: [https://github.com/vanessavcos/DFA-](https://github.com/vanessavcos/DFA-2D)
 189 2D.

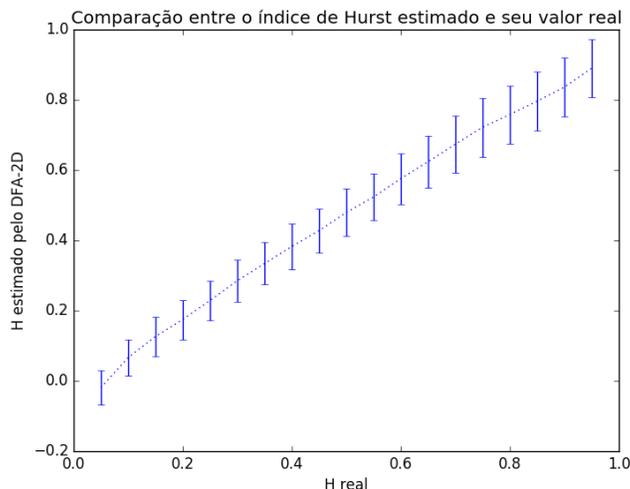


Figure 2: Mean value of the Hurst exponent estimated with the DFA-2D for fBm surfaces.

190 3.4 Process of Classification

191 Figure 3 and Table 3 show the log-log graph and the respective values of
 192 α for fGn and fBm with the Hurst exponent 0.5 and white ($\beta = 0$), pink
 193 ($\beta = 1$) and brown ($\beta = 2$) noises. The values are the averages of the
 194 five hundred samples of each model. The fGn data were produced by the
 195 successive differences of the respective fBm, in accordance with [36, 37].

196 From the value of Table 3, it was determined that the DFA-2D was able
 197 to discriminate between the different types of signals. It was also found that
 198 the α for the white and brownian noises have border values (approximately
 199 1 and 2, respectively) and reveal a close relationship with each other because
 200 the brownian noise is the integration of white noise. An important point is
 201 the evaluation of the DFA-2D for $(1/f)$ pink noises. The α is given a value
 202 of approximately 1.5, which reveals brownian signals for the DFA-1D. It
 203 should be noted that the DFA-2D returns values between 1 and 2 for noises

Signal	Parameter	α_{mean}
fGn	Hurst = 0.5	0.4067 ± 0.0205
fBm	Hurst = 0.5	2.4794 ± 0.0676
White Noise	$\beta = 0$	0.9992 ± 0.0207
Pink Noise	$\beta = 1$	1.4895 ± 0.0313
Brownian Noise	$\beta = 2$	1.9820 ± 0.0434

Table 3: Mean values of α to some typical signals obtained with DFA-2D. The mean and standard deviation are relative to five hundred samples of each model.

204 and that, in this case, there is no way to derive the Hurst exponent.

205 The fBm signal had a value very close to the expected (≈ 2.5). The fGn
 206 dispersed a little from the expected (≈ 0.5), which might be a reflection of
 207 the methodology employed to derive it from the fBm.

208 The log-log graph (Figure 1) makes clear the monofractality of the data
 209 because the linear regression is a straight line. A structural difference is also
 210 found between a fractionary gaussian noise (fGn) and an ordinary gaussian
 211 noise (white noise). The fGn shows the same value of the fluctuation function
 212 in all the scales because it is a stationary signal.

213 When the data from Table 1 and Table 3 are compared, it is found that
 214 the relation shown in eq. 6 cannot be confirmed. Instead of this, it can be
 215 said that for two-dimensional signals $\beta \equiv 2\alpha - 2$ and therefore, $\alpha \equiv \frac{\beta}{2} + 1$ for
 216 noises and fGn signals. In the case of fBm signals, it can be approximated
 217 that $\beta \equiv \alpha$. However, in the case of the fBm signal when the Hurst exponent
 218 = 0.5, the power spectrum overestimates the value of H and hence, the value
 219 of α , because the value obtained for β was 2.7, when it should have been
 220 equal to 2.5. Tests with other values of the Hurst exponent were carried out
 221 and these confirmed the relation of $\beta \equiv \alpha$ for the two-dimensional fBm.

222 3.5 Computational Performance

223 This section has evaluated the relation between the processing time, the
 224 accuracy of the DFA-2D and the size of the image. It involved assessing
 225 thirty samples of the fBm bi-dimensional process, with different sizes (64,
 226 128, 256, 512, 1024 e 2048) and different Hurst exponents (0.3, 0.5 e 0.8).
 227 The results are shown in Table 4.

228 With regard to accuracy, the images of 64x64 and 2048x2048 were those
 229 that had most errors. However, even in small-scale conditions, as in the case
 230 of the 64x64 image, the DFA-2D came close to the result expected.

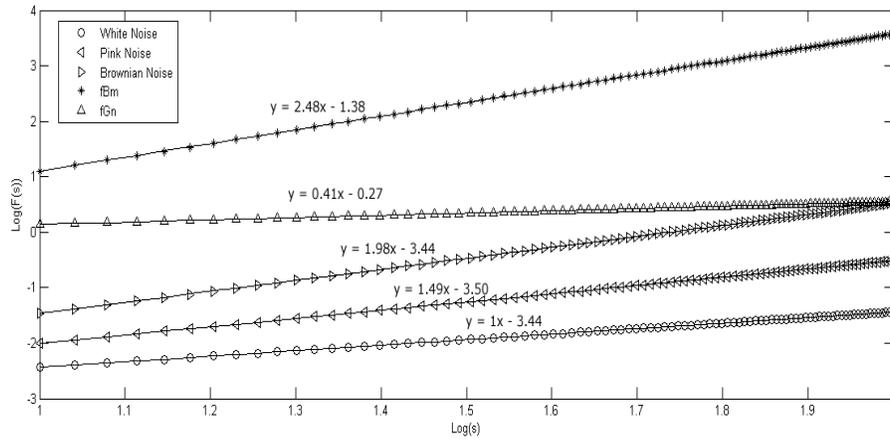


Figure 3: LogLog graphic of typical signals fBm and fGn with Hurst 0.5 and white, pink and Brownian noise. The data are averages of five hundred samples. The highlighted lines and equations refer to the fit of the data. The slope of this line reveals the value of the parameter α .

231 With regard to performance, it was confirmed in Table 4 and Figure
 232 4 that the enlargement of the image led to an exponential increase in the
 233 time. Some authors have drawn attention to the high computational cost
 234 of the DFA-2D (asymptotic complexity of the order of $O(n^3)$) [26, 27]. One
 235 way to improve performance might be by means of parallelization, since the
 236 sub-matrices can be processed in parallel because they do not depend on
 237 each other's results. The parallel algorithm that uses GPU and CUDA is
 238 under development and will be described shortly. The preliminary results
 239 suggest there was a significant improvement in the processing time when the
 240 parallel approach was adopted.

241 4 Concluding Remarks

242 The DFA is a method that is widely used in several areas. It is important
 243 to know the dynamics of data formation, especially in cases of classification.
 244 Despite some criticism, several variations of the method have been designed,
 245 which is proof of its robustness. However, the DFA-2D version has not been
 246 explored to any great extent and there are still very few evaluations with
 247 different data from fBm and fGn.

248 In this study, there is an evaluation of the DFA-2D for white, pink and

Hurst	Image Size	α		Time(s)	
		Mean	Std	Mean	Std
0.3	64	0.3218	0.0281	0.0911	0.0010
	128	0.3126	0.0129	0.4850	0.0031
	256	0.3030	0.0222	2.3569	0.0048
	512	0.3043	0.0126	11.6196	0.1045
	1024	0.3034	0.0094	59.9468	0.1172
	2048	0.2987	0.0545	357.1790	119.3266
0.5	64	0.5204	0.0403	0.0910	0.0013
	128	0.5037	0.0245	0.4837	0.0010
	256	0.5071	0.0256	2.3583	0.0127
	512	0.4954	0.0186	11.5791	0.0562
	1024	0.5033	0.0123	59.8843	0.0904
	2048	0.4989	0.0913	253.8633	146.4943
0.8	64	0.7892	0.0628	0.0911	0.0021
	128	0.8122	0.0445	0.4849	0.0028
	256	0.8054	0.0262	2.3574	0.0122
	512	0.8020	0.0271	11.5651	0.0176
	1024	0.8043	0.0268	59.8982	0.2121
	2048	0.7880	0.0150	356.1359	148.3523

Table 4: Mean and standard deviation of DFA-2D parameter α and the processing time for images with different sizes and different Hurst exponents.

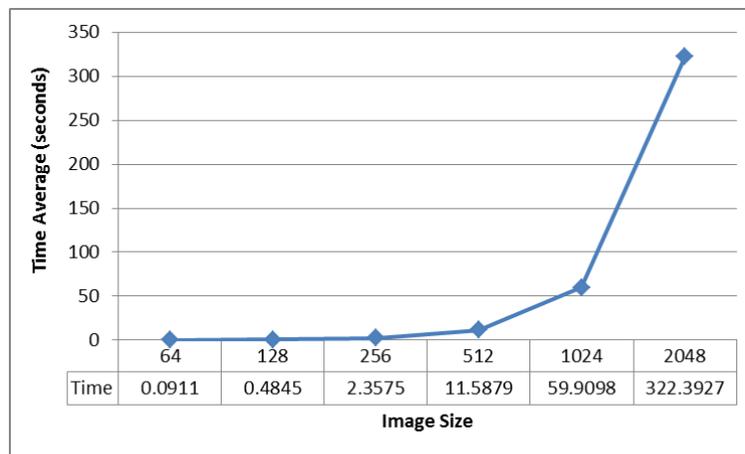


Figure 4: Average processing time values for different image sizes.

249 brownian noises. It was found that the classical relationship between α and
250 β is not valid for the data analyzed in this study where the exponents cited
251 were related in the form $\beta = 2\alpha - 2$, with the exception of the fBm signal,
252 where $\beta \equiv \alpha$.

253 The average time for the data processing (512 x 512) is about 12 seconds
254 for the DFA-2D, in a totally dedicated Linux machine. This high complex-
255 ity of the algorithm has led some sectors to adjust it or even avoid using
256 the method. For this reason, the parallelization of the DFA-2D have been
257 examined and the preliminary results from the use of GPU and CUDA have
258 shown a significant improvement in the results.

259 In the case of fBm signals, the accuracy of the method is not greatly
260 influenced by the image size. Even with small scales, such as in the image
261 of 64x64, the DFA-2D was able to come close to the result expected.

262 It was found that the DFA-2D was able to classify and characterize the
263 different systems studied and allow different signals to be identified. Another
264 valuable finding of the work is in the theoretical and practical understanding
265 of the white noise. It could be noted that, despite being decorrelated, the
266 two-dimensional signals with $\beta = 0$ (noise), $\alpha = 0.5$ (fGn) and $\alpha = 2.5$
267 (fBm), can physically identify different processes.

268
269 **Acknowledgments.** The authors would like to express their gratitude
270 to Professor Reinaldo Rosa (National Institute of Spatial Research Labo-
271 ratory for Computational and Applied Mathematics (LAC)) for his support
272 in this research undertaking.

273 References

- 274 [1] C. Peng, S. Buldyrev, A. Goldberger, S. Havlin, F. Sciortino, M. Si-
275 mons, and H. Stanley, "Long-range correlations in nucleotide se-
276 quences," *Phys. Rev. E*, vol. 49, pp. 1685–1689, 1994.
- 277 [2] J. W. Kantelhardt, S. A. Zschiegner, E. Koscielny-Bunde, S. Havlin,
278 A. Bunde, and H. E. Stanley, "Multifractal detrended fluctuation anal-
279 ysis of nonstationary time series," *Physica A*, vol. 316, no. 1-4, pp. 87–
280 114, 2002.
- 281 [3] G.-F. Gu and W.-x. Zhou, "Detrended fluctuation analysis for fractals
282 and multifractals in higher dimensions," *Physical Review*, vol. 74, no. 6,
283 pp. 1–7, 2006.

- 284 [4] H. Xiong and P. Shang, “Detrended fluctuation analysis of multivari-
285 ate time series,” *Communications in Nonlinear Science and Numerical*
286 *Simulation*, vol. 42, pp. 12–21, jan 2017.
- 287 [5] H. Hurst, “Long term storage capacity of reservoirs,” *Transactions of*
288 *thr American SOciety of Civil Engineers*, vol. 116, pp. 770–799, 1951.
- 289 [6] R. M. Bryce and K. B. Sprague, “Revisiting detrended fluctuation anal-
290 ysis,” *Scientific reports*, vol. 2, p. 315, jan 2012.
- 291 [7] Z. Chen, K. Hu, and H. E. Stanley, “Effect of nonstationarities on
292 detrended fluctuation analysis,” *Physical Review E*, vol. 65, no. 4, p. 15,
293 2002.
- 294 [8] M. Höll and H. Kantz, “The relationship between the detrendend fluc-
295 tuation analysis and the autocorrelation function of a signal,” *European*
296 *Physical Journal B*, vol. 88, no. 12, pp. 1–7, 2015.
- 297 [9] M. Höll and H. Kantz, “The fluctuation function of the detrended fluc-
298 tuation analysis - Investigation on the AR (1) process,” *European*
299 *Physical Journal B*, vol. 88, no. 126, pp. 1–15, 2015.
- 300 [10] J.-M. Lee, D.-J. Kim, I.-Y. Kim, K.-S. Park, and S. I. Kim, “Detrended
301 fluctuation analysis of EEG in sleep apnea using MIT/BIH polysomnog-
302 raphy data,” *Computers in Biology and Medicine*, vol. 32, pp. 37–47,
303 jan 2002.
- 304 [11] L. Parish, G. Worrell, S. Cranstoun, S. Stead, P. Pennell, and
305 B. Litt, “Long-range temporal correlations in epileptogenic and non-
306 epileptogenic human hippocampus,” *Neuroscience*, vol. 125, pp. 1069–
307 1076, jan 2004.
- 308 [12] S. L. Winter and J. H. Challis, “Classifying the variability in impact and
309 active peak vertical ground reaction forces during running using DFA
310 and ARFIMA models,” *Human Movement Science*, vol. 51, pp. 153–
311 160, jan 2017.
- 312 [13] R. G. Kavasseri and R. Nagarajan, “Evidence of crossover phenomena
313 in wind speed data,” *Circuits and Systems I: Regular Papers, IEEE*
314 *Transactions on*, vol. 51, no. 11, pp. 2255–2262, 2004.
- 315 [14] K. Koçak, “Examination of persistence properties of wind speed records
316 using detrended fluctuation analysis,” *Energy*, vol. 34, pp. 1980–1985,
317 nov 2009.

- 318 [15] S. Dutta, D. Ghosh, and S. Chatterjee, “Multifractal detrended fluctuation analysis of human gait diseases,” *Frontiers in Physiology*, vol. 4,
319 2013.
320
- 321 [16] A. Madanchi, M. Absalan, G. Lohmann, M. Anvari, and M. Reza
322 Rahimi Tabar, “Strong short-term non-linearity of solar irradiance fluctuations,” *Solar Energy*, vol. 144, pp. 1–9, mar 2017.
323
- 324 [17] E. Maiorino, L. Livi, A. Giuliani, A. Sadeghian, and A. Rizzi, “Multifractal characterization of protein contact networks,” *Physica A: Statistical Mechanics and its Applications*, vol. 428, pp. 302–313, jun 2015.
325
326
- 327 [18] A. V. Alpatov, S. P. Vikhrov, and N. V. Grishankina, “Revealing the surface interface correlations in a-Si:H films by 2D detrended fluctuation analysis,” *Semiconductors*, vol. 47, pp. 365–371, mar 2013.
328
329
- 330 [19] E. Barrera, F. Gonzalez, E. Rodriguez, and J. Alvarez-Ramirez, “Correlation of optical properties with the fractal microstructure of black molybdenum coatings,” *Applied Surface Science*, vol. 256, no. 6, pp. 1756–1763, 2010.
331
332
333
- 334 [20] Q. Nie, J. Xu, W. Man, and F. Sun, “Detrended fluctuation analysis of spatial patterns on urban impervious surface,” *Environmental Earth Sciences*, vol. 74, no. 3, pp. 2531–2538, 2015.
335
336
- 337 [21] C. Vargas-Olmos, J. S. Murguía, M. T. Ramírez-Torres, M. Mejía Carlos, H. C. Rosu, and H. González-Aguilar, “Two-dimensional DFA scaling analysis applied to encrypted images,” *International Journal of Modern Physics C*, vol. 26, no. 08, p. 1550093, 2015.
338
339
340
- 341 [22] O. Velazquez-Camilo, E. Bolaños-Reynoso, E. Rodriguez, and J. Alvarez-Ramirez, “Fractal analysis of crystallization slurry images,” *Journal of Crystal Growth*, vol. 312, no. 6, pp. 842–850, 2010.
342
343
- 344 [23] F. Wang, D.-w. Liao, J.-w. Li, and G.-p. Liao, “Two-dimensional multifractal detrended fluctuation analysis for plant identification,” *Plant Methods*, vol. 11, no. 12, pp. 1–11, 2015.
345
346
- 347 [24] R. P. Yadav, S. Dwivedi, A. K. Mittal, M. Kumar, and A. C. Pandey, “Fractal and multifractal analysis of LiF thin film surface,” *Applied Surface Science*, vol. 261, pp. 547–553, 2012.
348
349

- 350 [25] J. Alvarez-ramirez, E. Rodriguez, I. Cervantes, and J. Carlos, “Scal-
351 ing properties of image textures : A detrending fluctuation analysis
352 approach,” *Physica A*, vol. 361, pp. 677–698, 2006.
- 353 [26] A. J. EINSTEIN, H.-S. WU, and J. GIL, “Detrended Fluctuation Anal-
354 ysis of chromatin texture for diagnosis in breast cytology,” *Fractals*,
355 vol. 10, pp. 19–25, apr 2012.
- 356 [27] R. M. de Freitas, P. C. Genovez, M. C. Bentz, R. R. Rosa, and Y. E.
357 Shimabukuro, “Análise de anisotropia de imagens utilizando o método
358 DFA : um estudo de caso na área de exploração de petróleo,” in *Anais*
359 *XIV Simpósio Brasileiro de Sensoriamento Remoto Sensoriamento*,
360 pp. 6463–6470, 2009.
- 361 [28] R. F. Voss, “Evolution of long-range fractal correlations and $1/f$ noise
362 in DNA base sequences,” *Physical Review Letters*, vol. 68, no. 25,
363 pp. 3805–3808, 1992.
- 364 [29] Gildea, D. Gildea, T. Thornton, and M. Mallon, “ $1/F$ Noise in Human
365 Cognition,” 1995.
- 366 [30] J. M. Halley and W. E. Kunin, “Extinction Risk and the $1/f$ Family of
367 Noise Models,” *Theoretical population biology*, vol. 56, no. 3, pp. 215–
368 30, 1999.
- 369 [31] J. J. Lennon, “Red-shifts and red herrings in geographical ecology,”
370 *Ecography*, vol. 23, no. 1, pp. 101–113, 2000.
- 371 [32] G. Wornell, “Wavelet-based representations for the $1/f$ family of fractal
372 processes,” *Proceedings of the IEEE*, vol. 81, no. 10, pp. 1428–1450,
373 1993.
- 374 [33] J. Yearsley, “spatialPattern,” 2016.
- 375 [34] R. O. Prum, R. H. Torres, S. Williamson, and J. Dyck, “Coherent light
376 scattering by blue feather barbs,” *Nature*, vol. 396, pp. 28–29, nov 1998.
- 377 [35] E. Ruzanski, “Radially averaged power spectral density - Matlab code,”
378 2011.
- 379 [36] D. Delignieres, S. Ramdani, L. Lemoine, K. Torre, M. Fortes, and
380 G. Ninot, “Fractal analyses for ‘short’ time series: A re-assessment of
381 classical methods,” *Journal of Mathematical Psychology*, vol. 50, no. 6,
382 pp. 525–544, 2006.

- 383 [37] A. Eke, P. Hermán, J. B. Bassingthwaite, G. M. Raymond, D. B.
384 Percival, M. Cannon, I. Balla, and C. Ikrényi, “Physiological time series:
385 Distinguishing fractal noises from motions,” *Pflügers Archiv European*
386 *Journal of Physiology*, vol. 439, no. 4, pp. 403–415, 2000.
- 387 [38] B. B. Mandelbrot and J. W. Van Ness, “Fractional Brownian Motions,
388 Fractional Noises and Applications,” *SIAM Review*, vol. 10, pp. 422–
389 437, oct 1968.
- 390 [39] C. Heneghan and G. McDarby, “Establishing the relation between de-
391 trended fluctuation analysis and power spectral density analysis for
392 stochastic processes,” *Physical Review E*, vol. 62, no. 5, pp. 6103–6110,
393 2000.
- 394 [40] G. Heinzel, A. Rüdiger, and R. Schilling, “Spectrum and spectral den-
395 sity estimation by the Discrete Fourier transform (DFT), including a
396 comprehensive list of window functions and some new at-top windows,”
397 *Max Plank Institute*, pp. 1–84, 2002.
- 398 [41] S. V. Buldyrev, A. L. Goldberger, S. Havlin, R. N. Mantegna, M. E.
399 Matsa, C.-K. Peng, M. Simons, and H. E. Stanley, “Long-range corre-
400 lation properties of coding and noncoding DNA sequences: GenBank
401 analysis,” *Physical Review E*, vol. 51, pp. 5084–5091, may 1995.
- 402 [42] Scipy.org, “Numpy and SciPy Documentation,” 2016.
- 403 [43] INRIA, “FRACLAB, A Fractal Analysis Toolbox for Signal and Image
404 Processing,” 2016.